# U(N) Tools for Loop Quantum Gravity

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#### Started with F. Girelli,

then work with L. Freidel and with E. Borja, J. Diaz-Polo & I. Garay

arXiv:gr-qc/0501075, 0911.3553, 1005.2090, 1006.2451, 1006.5666, and more coming soon!

Etera Livine Ecole Normale Supérieure de Lyon - CNRS U(N) Tools for Loop Quantum Gravity

**Aim:** Use new U(N) tools to probe the structure of the intertwiner space of Loop Quantum Gravity and reformulate the LQG dynamics.

A U(N) Action on the Space of SU(2) Intertwiners

- Characterizes Intertwiner Spaces as U(N) Irreps
  - $\hookrightarrow$  Intertwiner counting and Black Hole entropy in LQG
- ② The Intertwiner Space as a  $L^2$  space with
  - $\hookrightarrow \mathsf{Creation}/\mathsf{annihilation}\ \mathsf{operators}$
  - $\hookrightarrow \mathsf{Semi-classical} \ \mathsf{coherent} \ \mathsf{intertwiner} \ \mathsf{states}$
- Improves the geometrical interpretation of Intertwiners
  - $\hookrightarrow$  Reconstructing spin networks from gluing classical poyhedra
- Provides new operators for the LQG Dynamics
  - $\hookrightarrow \mathsf{Reformulation} \text{ of Holonomy}/\mathsf{Flux} \text{ operators}$

The Loop Gravity Quantization Scheme

Loop Quantum Gravity = Canonical Quantization of GR

Based on a 3+1 formalism  $\mathcal{M} = \mathbb{R} \times \Sigma_{3d}$ 

- GR as a Gauge Theory :  $\{A_i^a(x), E_b^j(y)\} = \gamma \delta_i^j \delta_b^a \delta^{(3)}(x-y)$ 
  - $\hookrightarrow$  Triad field  $E_a^i$  gives 3-metric  $h^{ij} = E_a^i E_a^j$
  - $\hookrightarrow$  Ashtekar-Barbero  $\mathfrak{su}(2)$  connection  $A_i^a = \Gamma_i^a(E) + \gamma K_i^a$

 $\hookrightarrow \mathsf{Immirzi} \text{ parameter } \gamma$ 

- GR as a Constrained System :  $H = \Lambda^a \mathcal{G}_a + N_i \mathcal{H}^i + N \mathcal{H}$ 
  - $\hookrightarrow$  SU(2) gauge transformations  ${\cal G}_a \sim d_A E$
  - $\hookrightarrow$  3d space diffeomorphisms (on-shell) generated by  $\mathcal{H}^i \sim \textit{EF}[A]$
  - $\hookrightarrow \mathsf{Hamiltonian}\ \mathsf{constraint}\ \mathsf{encoding}\ \mathsf{time}\ \mathsf{evolution}\ \mathcal{H} \sim \mathsf{EEF}[A]$

Spin Network States and Intertwiners

A straightforward quantization procedure:

- Wave-functions  $\psi(A)$  with triad acting as  $\hat{E} \sim i\gamma \frac{\partial}{\partial A}$
- Choose SU(2)-inv cylindrical functions ψ({U<sub>e</sub>}<sub>e∈Γ</sub>) depending on holonomies along the edges of a graph Γ
- Identify the Spin Network basis of LQG quantum states:
   Label edges with spins j<sub>e</sub>, put intertwiners i<sub>v</sub> at vertices and glue them with the holonomies along the edges.

$$\psi_{j_e,i_v}^{(\Gamma)}(U_e) \equiv \operatorname{Tr} \otimes_v i_v \otimes_e D^{j_e}(U_e) \qquad j_1 \underbrace{j_e \qquad D^{j_e}(U_e)}_{j_n}$$

• Then implement  $\mathcal{H}^i, cH$  constraints on these states.

#### Intertwiners ?

At each vertex v, the intertwiner intertwines the representations living on the in-coming edges:

Intertwiner  $i_v = SU(2)$ -inv tensor in  $j_1^{(v)} \otimes .. \otimes j_n^{(v)} =$  singlet state

### Geometrical Interpretation

We build geometrical operators out of the triad  $\hat{E}$ . It provides spin networks with an interpretation as quantized discrete geometries.

- Elementary surface  $\mathcal{A}$  dual to edge e  $\Rightarrow \mathcal{A}^2 = \gamma^2 \vec{J}_e \cdot \vec{J}_e = \gamma^2 j_e(j_e + 1)$  easy  $\checkmark$ Only depends on the spins.
- Chunk of volume  $\mathcal{V}$  dual to vertex v, picks triplets of edges  $\Rightarrow \mathcal{V}^2 = \gamma^3 \vec{J}_{e_1} \cdot (\vec{J}_{e_2} \wedge \vec{J}_{e_3})$  harder... Depends on the intertwiners.

 $\label{eq:linear} \begin{array}{l} \mbox{Intertwiners} = \mbox{basic building block for quantum} \\ \mbox{geometry in LQG} \end{array}$ 

The U(N) structure of the intertwiner space

**The Object:** the space of intertwiners with *N* legs i.e of SU(2)-inv states in  $V^{j_1} \otimes .. \otimes V^{j_N}$  for arbitrary values of spins  $j_i \in \mathbb{N}/2$ .

**The first tool:** The Schwinger representation of  $\mathfrak{su}(2)$  in term of harmonic oscillators

We introduce  $a_i, b_i$  for each leg *i*.



$$[a_i,a_j^{\dagger}]=[b_i,b_j^{\dagger}]=\delta_{ij},$$

$$J_i^z = \frac{1}{2} (a_i^{\dagger} a_i - b_i^{\dagger} b_i) \rightarrow m_i$$
$$J_i^+ = a_i^{\dagger} b_i, \ J_i^- = a_i b_i^{\dagger},$$
$$E_i = (a_i^{\dagger} a_i + b_i^{\dagger} b_i) \rightarrow j_i$$

The U(N) structure of the intertwiner space

We build SU(2)-invariant quadratic operators acting on pairs of legs:

$$E_{ij}\equiv (a_i^{\dagger}a_j+b_i^{\dagger}b_j)$$

They shift the spins  $j_i$  and form a  $\mathfrak{u}(N)$ -algebra :

$$[E_{ij}, E_{kl}] = \delta_{jk} E_{il} - \delta_{il} E_{kj}$$

The diagonal components give the spins living on the edges:

$$E_{ii} = E_i = \text{ spin } 2j_i,$$
  

$$E = \sum E_i = 2 \times \sum_i j_i = 2 \times \text{ total area,}$$
  

$$[E, E_{ij}] = 0.$$

# The U(N)-action on Intertwiners

The basic conclusion: unitary group U(N) acts on intertwiners.

After a little more work, U(N) acts on the space of intertwiners with fixed N (obviously) and with fixed total boundary area (U(1) Casimir E):

$$R_N^J = \bigoplus_{\sum_{i=J}^N j_i = J} \operatorname{Inv}[V^{j_1} \otimes .. \otimes V^{j_N}]$$

Irreducible representations with highest weight vector satisfying:

 $E_1|v\rangle=E_2|v\rangle=J|v\rangle, \quad E_{k\geq 3}|v\rangle=0 \ \Rightarrow \ {\rm bivalent \ intertwiner}.$ 

The operators  $E_{ij}$  shift the spins on the legs and allow to go from this highest weight vector to arbitrary intertwiner states.

# More SU(2)-invariant operators !

We identify a new set of SU(2)-invariant quadratic operators:

$$F_{ij} \equiv a_i b_j - a_j b_i, \quad F_{ij} = -F_{ji}.$$

They form a closed algebra with the  $E_{ij}$ . They are invariant under SU(2) but do not preserve the total area!!

$$[E, F_{ij}] = -2F_{ij}, \qquad [E, F_{ij}^{\dagger}] = +2F_{ij}^{\dagger}$$

# $\Rightarrow$ annihilation/creation operators

They allow to shift between U(N) irreps and provide the space of intertwiners with a Fock space interpretation:

$$R_N = \bigoplus_{\{j_i\}} \operatorname{Inv}[V^{j_1} \otimes .. \otimes V^{j_N}] = \bigoplus_{J \in \mathbb{N}} R_N^J.$$

What has been done with the U(N) framework?

- The original objective: We replaced the scalar product operators J<sub>i</sub> · J<sub>j</sub> by another set of SU(2)-invariant operators E<sub>ij</sub>, which form a closed u(N) algebra.
   → Useful to define semi-classical states.
- U(N) = unitary deformations of intertwiners at fixed area.
  - $\hookrightarrow$  Useful to define evolution for LQG?
- Intertwiner Counting from the dimensions of U(N) irreps
   → Applied to LQG black hole entropy calculations
- Constructing coherent intertwiners using the creation operators (F<sup>†</sup>)<sup>J</sup>
  - $\hookrightarrow$  Coherent states under  $\mathrm{U}(N)$  transformations
  - $\hookrightarrow$  Peak the values of scalar products
  - $\hookrightarrow$  Useful to define semi-classical spin network states

# Going further with U(N) tools: Two developments

• The intertwiner space as a  $L^2$  space :

$$R_{N} = L^{2}_{holo} \left( \frac{\mathrm{U}(N)}{\mathrm{U}(N-2) \times \mathrm{SU}(2)} \right) = L^{2}_{holo} \left( \frac{\mathbb{C}^{2N}}{\mathbb{C} \times \mathrm{GL}(2,\mathbb{C})} \right)$$

 $\rightsquigarrow$  as wave-functions on U(N) or of N spinors.

 $\Rightarrow$  Allows to truly view intertwiners as quantized polyhedra.

 Application to the 2-vertex model : Truncation of QG to spin networks living on the 2-vertex graph (proposed by Rovelli & Vidotto) → Established a dictionary between U(N)-operators and holonomy /grasping operators of LQG → Use the U(N) symmetry to average over all intertwiners with same total boundary area, and effectively reduce to isotropic states i.e impose spherical symmetry . ⇒ Reducing from LQG to (L)QC at the quantum level... Where are we going with the U(N) framework for Loop Quantum Gravity?

Some direction for the future (of the U(N) formalism):

- A reformulation of LQG dynamics in terms of spinors and U(N) operators acting directly on the vertices/intertwiners of the spin network states?
- A better understanding of how to write semi-classical states for the quantum geometry?
- Progress on understanding how to deform spin network states using U(N) transformations →→ discrete diffeomorphisms?
- Implement the symmetry reduction to quantum Cosmology (and possibly other symmetry reduced settings in general relativity) directly at the quantum level from Loop Quantum Gravity → an explicit bridge between LQG and LQC?
- And now to lñaki 's talk!