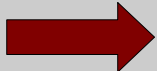


# Non-linear preheating after inflation and gravitational wave production

- Inflation  (p)reheating
- Preheating:
  - Field + metric perturbations
  - Preheating of tensor modes: Hybrid models
- Summary

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MBG, M. Tristram,  
J. Macias-Perez, D. Santos PRD77 '08  
MBG, J. Macías-Pérez, D. Santos, PRL105 '10

# Expanding Universe

Flatness problem

$$\Omega_T = 1 \quad \rightarrow \quad \Omega_T(t_{\text{nucl}}) - 1 \approx 10^{-16}$$

Horizon problem

The observable Universe was larger than the **particle horizon** at LSS

**Inflation**

Early period of accelerated expansion

$$\ddot{a} > 0: \quad P < -\rho/3$$

Super-horizon perturbations?

Too small sub-horizon  
(**causal**) perturbations

Unwanted relics

**monopoles**, moduli, gravitinos,...

Starobinsky '80; Guth '81; Albrecht, Steinhardt '82; Linde '82

# Slow Roll Inflation

Scalar field rolling down its (flat) potential

- Quantum fluctuations  $\rightarrow$  Primordial Spectrum

$$P_R^{1/2} \simeq (H_k / \dot{\phi}) P_{\delta\phi}^{1/2} \simeq 5 \times 10^{-5} \quad (k \simeq a_0 H_0) \quad \text{(Scalar)}$$

**(Superhorizon)**

- Primordial Gravitational Waves:

$$ds^2 = dt^2 - a(t)^2 (\delta_{ij} (1 + 2\psi) + h_{ij}) dx^i dx^j \quad P_T^{1/2} \sim 2\sqrt{2}H / (2\pi m_p)$$

**Tensor-to-scalar ratio  $r < 1$**



Oscillations : Matter  
Radiation

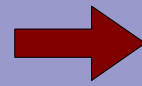
$$\dot{\rho}_\phi + 3H\rho_\phi + \Gamma_\phi\rho_\phi = 0$$

$$\dot{\rho}_R + 4H\rho_R - \Gamma_\phi\rho_\phi = 0$$

The inflaton interacts with other fields

# Preheating

$$V = V(\phi) + \frac{g^2}{2} \phi^2 \chi^2$$



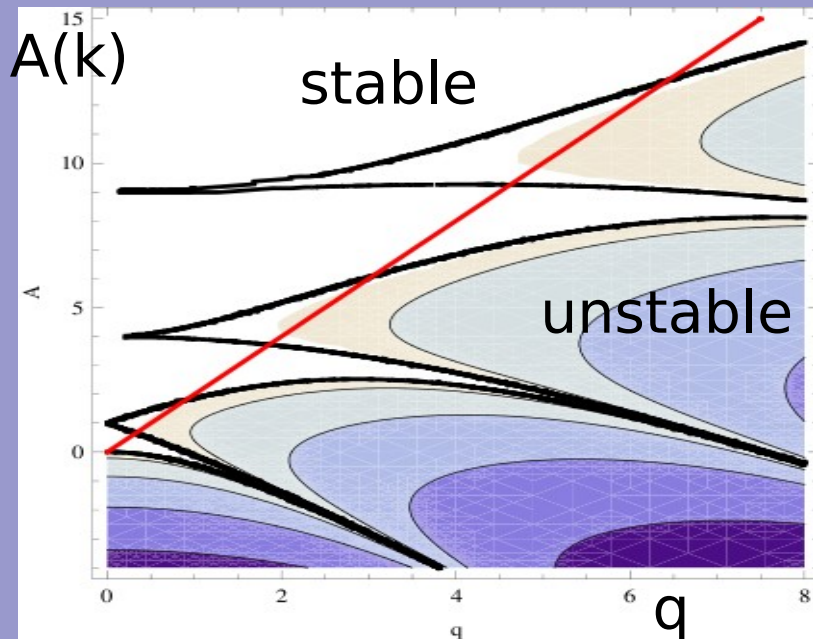
$$m_\chi^2(t) = g^2 \phi^2(t)$$

- Non-adiabatic change in the time dependent effective masses leads to parametric amplification of the field fluctuations

$$\delta \ddot{\chi}_k + 3H \delta \dot{\chi}_k + (k^2/a^2 + m_\chi^2(t)) \delta \chi_k = 0$$

$$\dot{m}_\chi(t)/m_\chi^2(t) \gg 1 \rightarrow \delta \chi_k \sim \exp(2\mu_k \omega t)$$

Particle production within certain resonance bands in k-space



- Mathieu equation:

$$\psi_k'' + (A(k) - 2q \cos 2z) \psi = 0$$

$$(\psi_k = a^{3/2} \delta \phi, z = \omega t)$$

$$A(k) = 2q + k^2/\omega^2 a^2$$

$$q = \langle m^2(t) \rangle / 4\omega^2$$

$q < 1$  : narrow resonance

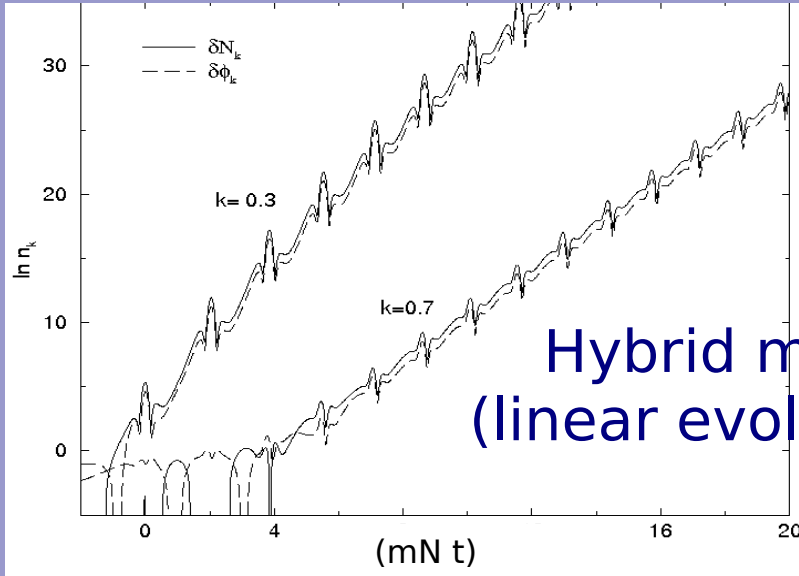
$q \gg 1$  : broad resonance

Dolgov & Kirilova '90;

Traschen & Brandenberger '91; Kofman, Linde, Starobinsky '94, '97

# Preheating

- Linear perturbation theory breaks down when



$$\langle (\nabla \chi)^2 \rangle \sim \langle \dot{\phi}^2 \rangle \sim V(\phi)$$

Gradient  $\sim$  Kinetic  $\sim$  Potential

$$\ddot{\chi}(\mathbf{x}) + 3H\dot{\chi}(\mathbf{x}) - \frac{1}{a^2} \nabla^2 \chi(\mathbf{x}) + \underline{g^2 \chi(\mathbf{x}) \phi(\mathbf{x})^2} = 0$$

Mode-to-mode coupling



backreaction + rescattering

- Non-perturbative process : numerical tools (Lattice)

IR cutoff  $k_{\min} \sim (2\pi)/L$

UV cutoff  $k_{\max} \sim N(2\pi)/L$

Parametric resonance  $\longrightarrow$  turbulence  $\longrightarrow$  thermalization  
 (Non thermal spectrum)  $\langle \chi^2 \rangle \propto t^{-2n}$

Khlebnikov & Tkachev '96; Prokopec & Roos '97;  
 Felder et al. '01; Micha & Tkachev '02

- Einstein Equations:  
non-linear coupled system matter - metric

Preheating of field fluctuations may lead to:

- Amplification of super-Hubble perturbations  
(primordial spectrum) Tsujikawa & Bassett '02
- Non-gaussianity Enqvist et al. '05 ; Barnaby & Cline '06  
Chambers & Rajantie '08
- Stochastic background of gravitational waves  
Easter, Giblin & Lim PRL99 '07; PRD77 '08  
García-Bellido & Figueroa PRL98 '07  
García-Bellido, Figueroa & Sastre PRD77 '08  
Dufaux et al. PRD76 '07; JCAP 03 '09  
MBG, Macías-Pérez, Santos PRL105 '10

Previous studies  $ds^2 = dt^2 - a(t)^2 (\delta_{ij} (1 + 2\Psi) + h_{ij}) dx^i dx^j$

$$\ddot{h}_{ij} + 3H \dot{h}_{ij} - \nabla^2 h_{ij} / a^2 = S_{ij}^{TT} / a^2 \quad \text{transverse \& traceless source term}$$

$$\underline{S_{ij}^{(\Psi)}} \ll S_{ij}^M \quad \text{Only matter (fields)}$$

- Fields evolve in a background metric
- Easter & Giblin & Lim

$$\text{source } S_{ij}^M(t, \mathbf{x}) \sim 16\pi G \nabla_i \chi \nabla_j \chi \quad \xrightarrow{\text{FT}} \quad S_{ij}^{TT}(t, \mathbf{k}) = \Lambda_{ij,lm}(\mathbf{k}) S_{ij}(t, \mathbf{k})$$

Tensors evolve in momentum space

- Dufaux & Bergman & Felder & Kofman & Uzan

$$\text{Green function method: } h_{ij}(t, \mathbf{k}) = 16\pi G \int G(t, t') S_{ij}^{TT}(t', \mathbf{k}) / k dt'$$

- García-Bellido & Figueroa & Sastre

$$\text{Tensors evolve in a lattice with } S_{ij} \quad h_{ij}^{TT}(t, \mathbf{k}) = \Lambda_{ij,lm}(\mathbf{k}) h_{ij}(t, \mathbf{k})$$

# Fields + metric perturbations in a lattice

Einstein eq.:  $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = T_{\mu\nu}/m_P^2$  Matter fields  $\Phi(x,t)$

- BSSN formalism + synchronous gauge

$$ds^2 = dt^2 - e^{4\beta(t,x)} \tilde{\gamma}_{ij} dx^i dx^j \quad \text{Det}(\tilde{\gamma}_{ij}) = 1$$

$\langle e^{2\beta(t,x)} \rangle \sim a(t)$  scale factor       $\langle 2\dot{\beta}(t,x) \rangle \sim H(t)$  Hubble rate

$\tilde{\gamma}_{ij}(t,x) \sim \delta_{ij} + h_{ij}(t,x) \sim$  tensors

- EOM:  $\ddot{\beta} + 2\dot{\beta}^2 = -(\rho + 3P)/(3m_P^2) - \dot{\tilde{\gamma}}_{ij}\dot{\tilde{\gamma}}^{ij}/24$

$$\dot{\tilde{\gamma}}_{ij} + 6\dot{\beta}\tilde{\gamma}_{ij} = 2e^{-4\beta}(M_{ij}^{\text{TF}} - R_{ij}^{\text{TF}}) + \dot{\tilde{\gamma}}_{il}\dot{\tilde{\gamma}}^l_j \quad \text{Traceless}$$

**fields                  metric**

➔ Fields ➔ scalar  $\beta$ :  $\underline{M_{ij}^{\text{TF}} \simeq R_{ij}^{\text{TF}}} \left\{ \begin{array}{l} M_{ij}^{\text{TF}} \simeq m_P^{-2}(\partial_i\Phi\partial_j\Phi)^{\text{TF}} \\ R_{ij}^{\text{TF}} \simeq (-4\partial_i\beta\partial_j\beta + 2\partial_i\partial_j\beta)^{\text{TF}} \end{array} \right.$

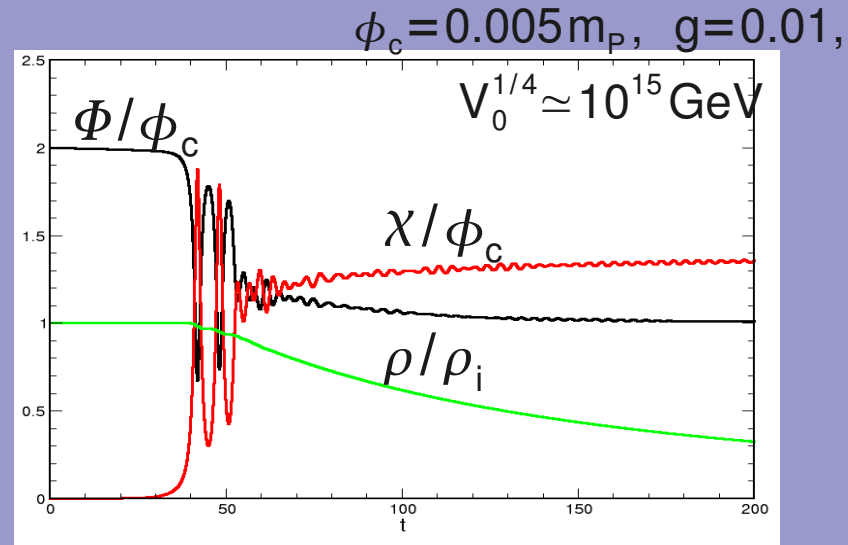
➔  $\dot{\tilde{\gamma}}_{ij}^{\text{TT}}(t, \mathbf{k}) = \Lambda_{ij,lm}(\mathbf{k}) \dot{\tilde{\gamma}}_{ij}(t, \mathbf{k})$  **Transverse & Traceless**



# Fields + metric perturbations: Hybrid model

$$V = V_0 + \frac{g^2}{4} \chi^4 + g^2 (\Phi^2 - \phi_c^2) \chi^2 + \frac{1}{2} m_\phi^2 \Phi^2$$

$$\Phi < \phi_c: \quad m_\chi^2(\Phi) < 0$$



Tachyonic preheating: very efficient particle production mechanism  $\delta \phi_k \simeq \delta \chi_k$

## Initial conditions:

No tensors at  $t=0$ ,

local scale factor and Hubble rate from constraint equations

Fields: slow-roll initial background values + vacuum fluct.

GW energy density:  $\rho_{\text{GW}} = m_{\text{P}}^2 \langle \dot{\gamma}_{ij}^{\text{TT}} \dot{\gamma}^{\text{ijTT}} \rangle / 4$

**Power spectrum:**  $\frac{d\Omega_{\text{GW}}}{d \ln k} = \frac{1}{\rho_c} \frac{d\rho_{\text{GW}}}{d \ln k} = \frac{\pi k^3}{3H^2 L^2} |\dot{\gamma}_{ij}^{\text{TT}}(k) \dot{\gamma}^{\text{ijTT}}(k)|$

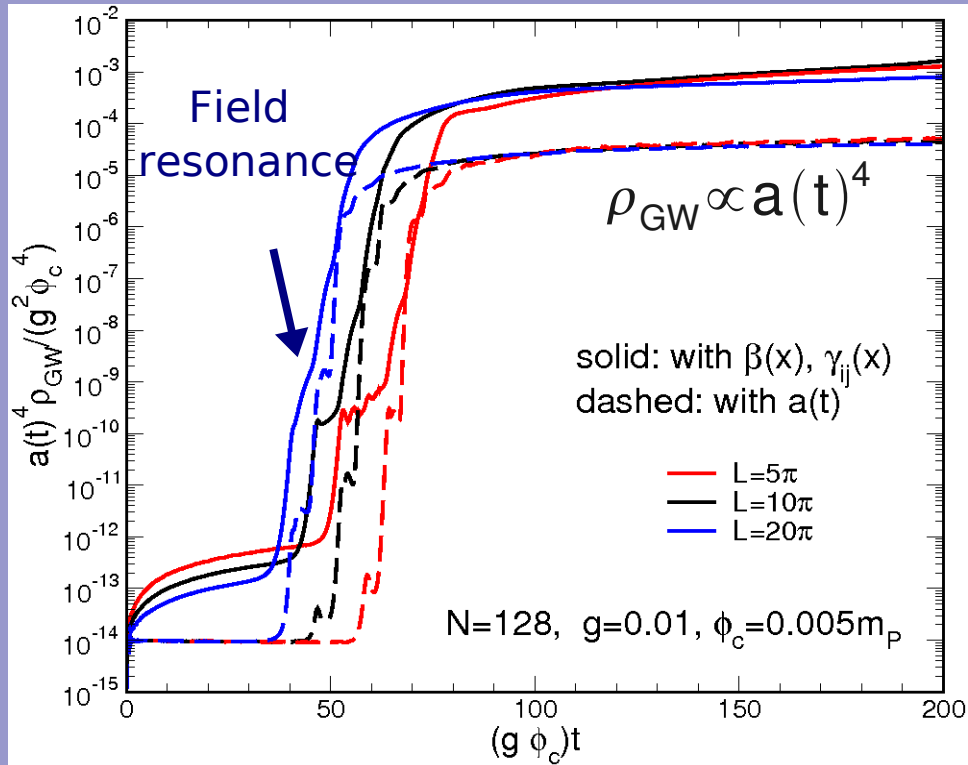


**Peak at  $k_{\text{res}} \sim g \Phi_c$**

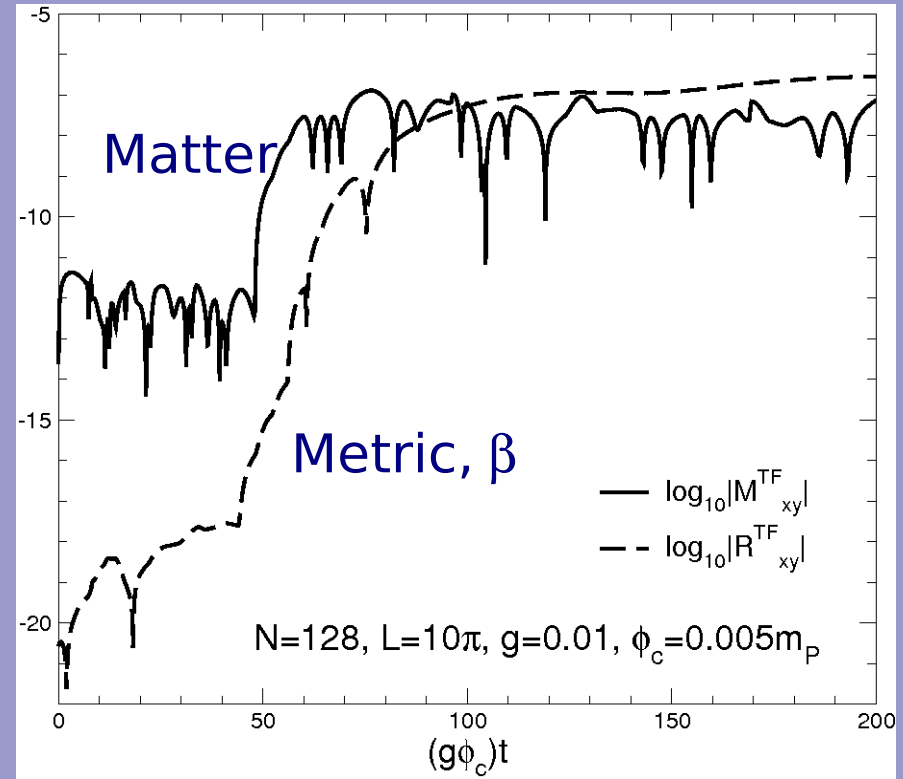
$\rho_{\text{GW}} \sim H_{\text{res}} / m_{\text{P}} \sim g \Phi_c^2 / m_{\text{P}}^2$

# Hybrid model

## GW energy density

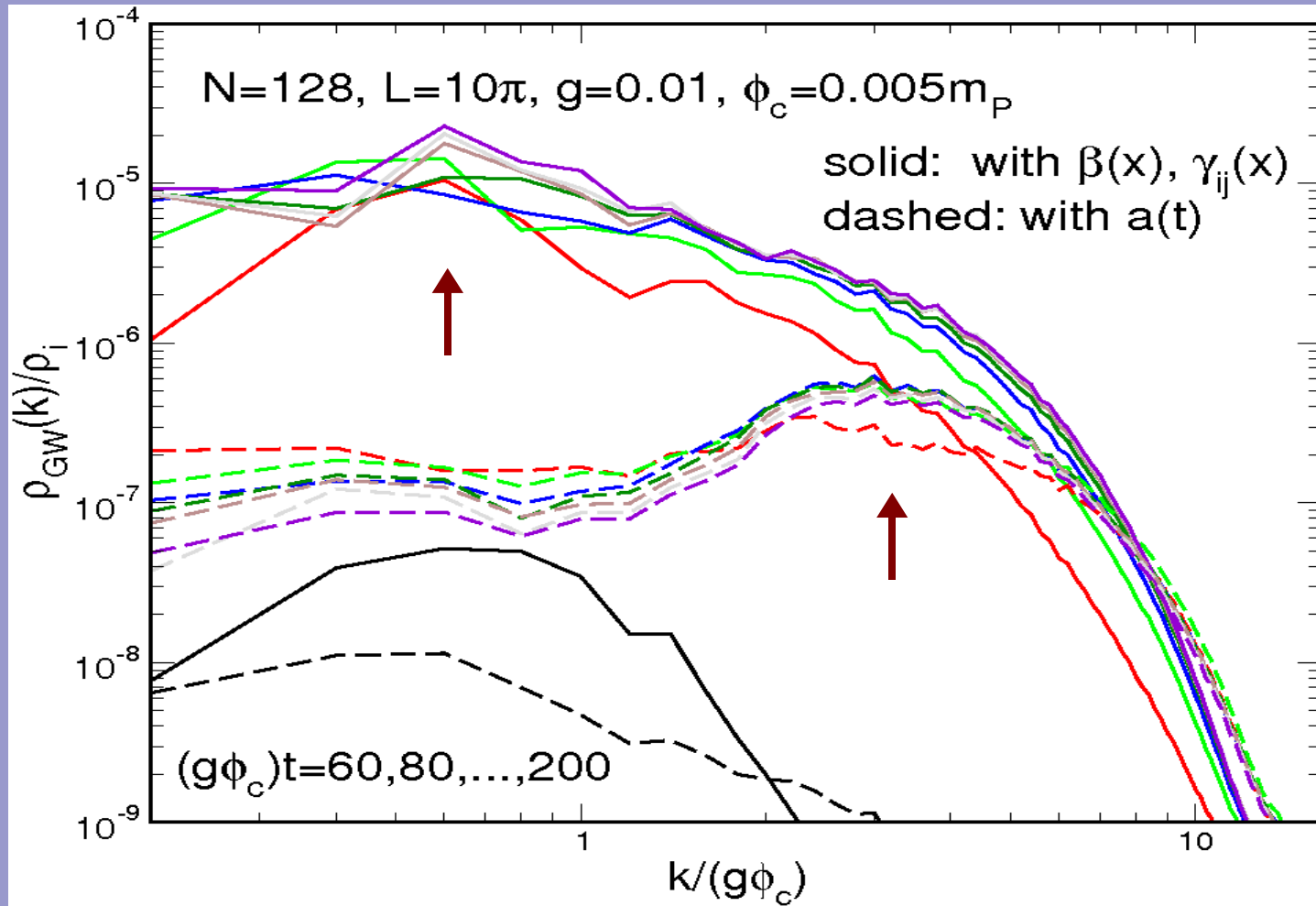


## Source terms



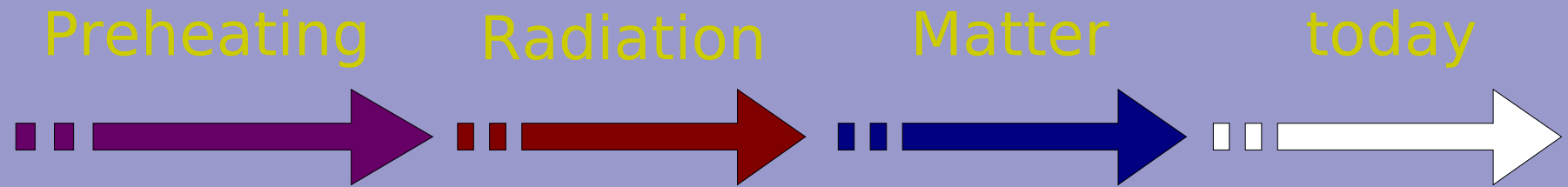
- larger  $\rho_{\text{GW}}$  ( $\sim$  an order of magnitude) when including metric perturbations

# Hybrid model : power spectrum



- $k_{\text{res}} \sim (g \Phi_c)$ , shifted towards smaller values
  - larger lattice needed to improve infrared resolution with a reliable ultraviolet cutoff
- $(k_{\text{min}}=0.4g \Phi_c)$

- From preheating to today values



**No entropy production**

$$g_S(T_R) a^3(t_R) T_R^3 = g_S(T_0) a^3(t_0) T_0^3$$

**Frequency:**

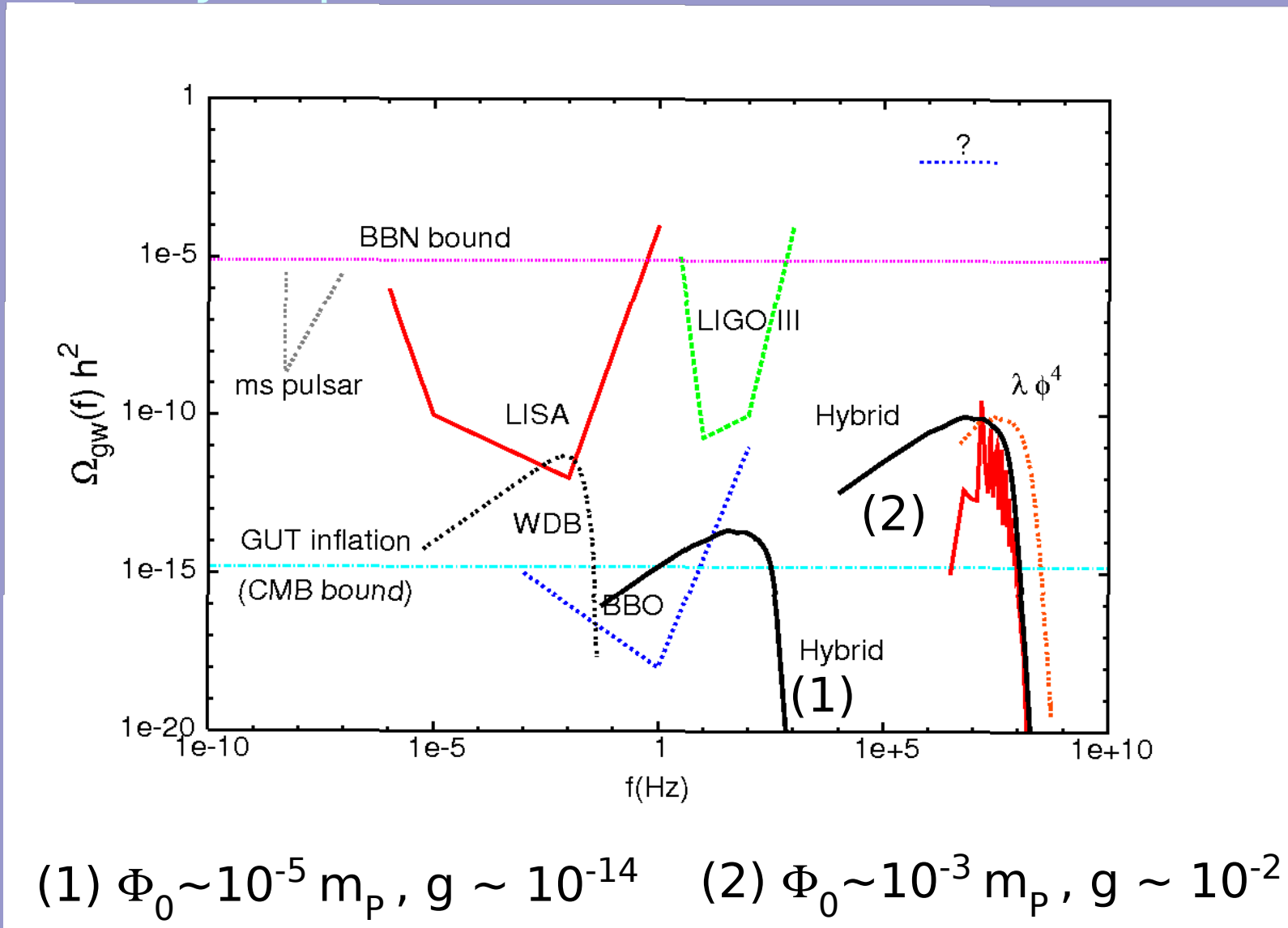
$$f_0 = \frac{k_0}{2\pi} \simeq \frac{k_{\text{res}}}{2\pi} \frac{a_{\text{res}}}{a_0} \simeq 6.4 \times 10^{10} \sqrt{g} \frac{k_{\text{res}}}{\rho_{\text{res}}^{1/4}} \text{ Hz} \sim 10^{10} \sqrt{g} \text{ Hz}$$

**Power spectrum:**

$$\Omega_{\text{GW}} h^2 = \Omega_{\text{GW}}^{(\text{res})} h^2 \left( \frac{a_{\text{res}}}{a_0} \right)^4 \simeq 5.5 \times 10^{-9} \left( \frac{\phi_c / m_{\text{P}}}{0.005} \right)^2$$

Low scale hybrid inflation (very weak coupling  $g < 10^{-14}$ ) could lead to a stochastic background of GW within the reach of GW detectors  $f_{\text{detec}} < 10^3 \text{ Hz}$

# Sensitivity of planned GW interferometers detectors




Very weak coupling  $g \ll 1$   $\longrightarrow$  tensor-to-scalar  $r \ll 1$

# Summary

- Preheating: parametric amplification of field + metric fluctuations during oscillations after inflation

$$\ddot{\Phi} + 6\dot{\beta}\dot{\Phi} - e^{-6\beta}\nabla^2\Phi + V_\phi - 2e^{-4\beta}\nabla^i\beta\nabla_i\Phi = 0$$


$$\ddot{\beta} + 2\dot{\beta}^2 \simeq -(\rho + 3P)/(3m_P^2)$$

$$\ddot{h}_{ij} + 3H\dot{h}_{ij} - e^{-2\beta}\nabla^2 h_{ij} \simeq S_{ij}^\Phi/m_P^2 + S_{ij}^\beta$$

$$|\beta_k|^2 \sim (\phi_c/m_P)^2 |\phi_k|^2$$

Scalar metric perturbations follow the same resonance pattern than the fields

Tensors sourced by 2<sup>nd</sup> order

matter and metric perturbations:  $S_{ij}^\beta \sim S_{ij}^\Phi/m_P^2$

- Coupled system fields + (scalar, tensor) metric fluct:
  - ~ 1 order of magnitude enhancement for  $\rho_{GW}$
  - ~ frequency shifted to smaller values

- Next step:

improve infrared and ultraviolet limits

Vectors? Magnetic fields?

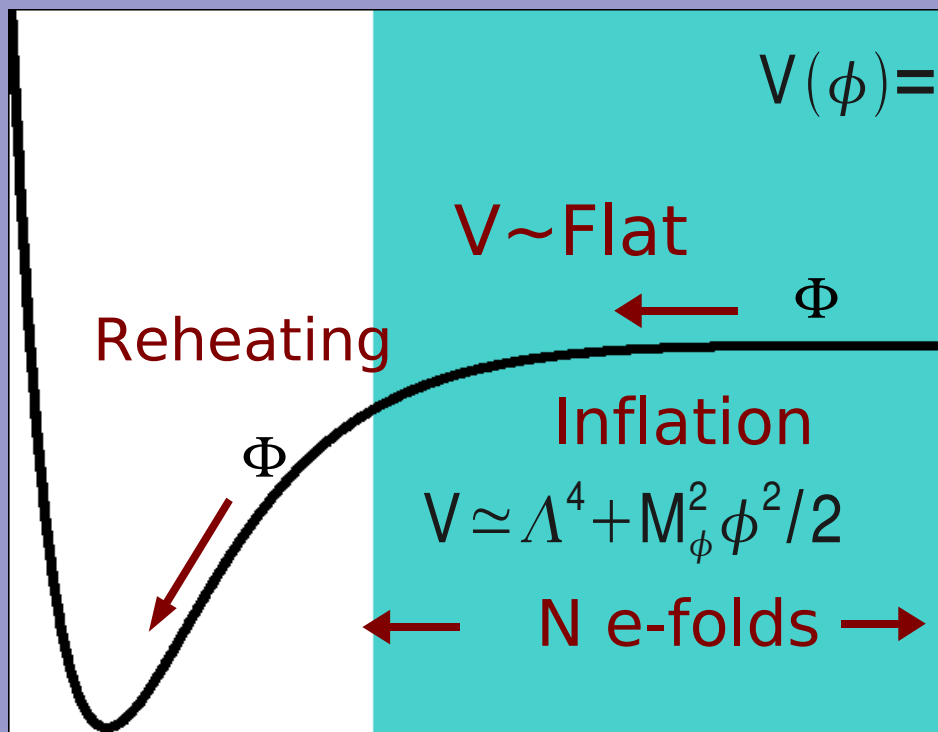
Non-gaussianity?

# Hybrid inflation (2 fields models)

Inflation  $\Rightarrow$  phase transition  $\Rightarrow$  reheating

False vacuum

Global minimum



$$V(\phi) = \Lambda^4 + \frac{M_\phi^2}{2} \phi^2 + \frac{\lambda}{4} \chi^4 + \frac{1}{2} (g^2 \phi^2 - \sqrt{\lambda} \Lambda^2) \chi^2$$

**Inflation:** above the critical value

The second field vanishes  
Potential  $\sim$  vacuum energy

**Global minimum:** below critical

The second field gets a vev  
The potential vanishes

Small (below Planck) field values  
**(No tensor:  $r \ll 1$ )**

# Cosmic Microwave Background Radiation

## Wilkinson Microwave Anisotropy Probe

